# Exploration of cyclopropyl radical ring opening to allyl radical by Newton trajectories: importance of valley-ridge inflection points to understand the topography 

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#### Abstract

Valley-ridge inflection points (VRI) on the potential energy surface for the ring opening of the cyclopropyl radical to allyl radical are determined using the tool of Newton trajectories (Quapp and Schmidt in Theor Chem Acc 128:47, 2011). VRIs play a role in the understanding of bifurcating reactions. The region where the bifurcation takes place is usually governed by a VRI point. For this important ring opening, the knowledge of the VRI point after the transition state was demanded some years ago (Mann and Hase in J Am Chem Soc 124:3208, 2002). Because the transition state is not symmetric, also the steepest descent from the transition state is not along a symmetry axis, and in such cases the steepest descent can fail the VRI point. That is the case here, indeed, though the pathway of the steepest descent goes near to a VRI point


[^0]downhill. However, an electronic intersection seam disturbs the relations. The exploration of the notorious curvilinear potential energy surface of this ring opening has delivered some further VRI points, which are reported. They give a frame for possible ring-opening channels including the conrotatory case.

Keywords Valley-ridge inflection point • Newton trajectory • Reaction bifurcation • Ring opening of cyclopropyl radical • Disrotatory and conrotatory case

## 1 Introduction: the theory of Newton trajectories and bifurcation points

The concepts of potential energy surface (PES), reaction path (RP), and its more restrictive definition minimum energy path (MEP) are the basic grounds of many theoretical chemistry theories and models [1, 2]. The RP is defined as a continuous curve in coordinate space, which connects two minimums of the PES by passing through a first order saddle point (SP), also called transition structure (TS), of the PES. The energy of the SP is assumed to be the highest value tracing along the RP. The chemical kinetic theories are based implicitly or explicitly on this definition. Many different curves satisfy the RP definition. This is the reason of a variety of curves used for this purpose. The most common used are the steepest descent (SD) or gradient path [3], gradient extremals (GE) [4-6], reduced gradient following (RGF) [7-14] and its equivalent definition, the so-called Newton trajectories (NT) [15]. If the NT is completely located in a valley region, then this RP meets the category of MEP. That a curve type used as an RP representation becomes MEP is possible for different mathematical definitions.

An important feature of the PES is that a valley or a ridge can bifurcate. The fact is related to the existence of valley-ridged inflection (VRI) points [15]. A VRI point is a point in the configuration space where, orthogonally to the gradient, at least one main curvature of the PES becomes zero. The gradient itself is assumed to be not the zero vector. If additionally the gradient is zero, then a stationary point and a VRI point coincide [16]. It is important to take into account that VRI points, which are defined in this way, are given independently of a curve definition of the RP representation. The VRI points are features of the PES under consideration. On the other side, a branching point or bifurcation point (BP) of a curve selected as a representation of an RP is a point where the curve branches. Usually such points do not coincide with a VRI point except for special NTs. The NTs branching at a VRI point are called singular NTs. The great remainder are regular NTs, which also do not bifurcate.

The VRI points can be classified into two main classes. A valley starting from an SP can bifurcate downhill and the two branches can lead to two valleys with their corresponding minimums. Between the two valleys a ridge emerges. Then, the VRI point is a valley-pitchfork (vpVRI) bifurcation. Or, a valley starting from a minimum can bifurcate uphill and the two branches can lead to two SPs. Between the two valleys, again a ridge emerges leading usually to an SP of index two. Then, the VRI point is also a vpVRI bifurcation. There is another possibility that a ridge on the PES bifurcates into two ridges going uphill, and between the two ridges emerges (in the chemical "interesting" case) a valley. The VRI point is a ridge-pitchfork (rpVRI) bifurcation. A possibility, which is not as interesting from the point of view of chemistry, is that in between a ridge of index 2 emerges-in cases with a PES of more than two dimensions. A further unusual possibility is that the VRI point is of a valley-ridge touching kind. The last possibility is the border-line case due to flat branches. So to say, its character is of a mixed type [15, 17]. The possibilities, vpVRI and rpVRI (interesting and may be not so interesting cases of higher index), as well as the mixed type, will emerge on the PES of the cyclopropyl radical ring opening, which is analyzed throughout the article.

Between a minimum and an SP of index 1, there are infinitely many regular NTs. However, between a minimum and an SP of index 2 is one singular NT that also crosses the VRI point in between. In Ref. [18] an iteration is proposed to find exactly the one singular NT using an empirical variational approach. We use the method throughout this article. With an initial choice of a search direction, we calculate by predictor and corrector steps an NT starting from a minimum, or an SP into the region of the guessed VRI. NTs are continuous curves which we represent by nodes in certain distances. The steplength of
the predictor step prescribes the distance. Connecting all nodes of the NT with the guessed VRI and dividing the lines into further nodes, we test for all nodes the value of norm $|\mathbf{A g}|$ and use the node with the minimal value for the new VRI point. Note that at the VRI point holds $|\mathbf{A g}|=0$. $\mathbf{A}$ is the adjoint matrix to the Hessian, and $\mathbf{g}$ is the gradient of the PES. The gradient at the guessed VRI is then used for the new search direction for a new NT run, and so on, up to convergence. Of course, the key always is an appropriate guess of the VRI region, at the beginning. The method is coupled with the GAMESS-US program ${ }^{1}$ for the calculation of energy, gradient, Hessian matrix, and metric matrices in non-redundant, internal z-matrix coordinates. All this is used at every node without any updating, see also the web-page [21] for the programs on the VRI-propose [18].

## 2 Thermal ring opening of cyclopropyl radical into allyl radical

In Organic Chemistry and Chemistry, in general the prediction of stereochemistry of the electrocyclic ring-opening reactions, such as cyclopropyl radical, has been a longstanding question. For cyclic molecular structures in their ground state with an even number of electrons, these reactions are governed by the Woodward-Hofmann rules [22-24]. These rules predict either a conrotatory or disrotatory stereochemistry evolution depending on the orbital diagram associated with the system. For the cyclopropyl radical, a system with an odd number of electrons, the Woodward-Hofmann rules predict that both the conrotatory and the disrotatory stereochemistry evolution is nominally forbidden [25]. In the computational study reported in Ref. [25], a highly asynchronous transition structure with $\mathrm{C}_{1}$ symmetry was identified. In a later computational study at the B3LYP/6-311G(2d) level of theory [26], the calculations of the RP represented as Intrinsic Reaction Coordinate (IRC) were carried out from the $\mathrm{C}_{1}$ transition state to the allyl radical, concluding that the overall reaction occurs with disrotatory stereochemistry. Nevertheless, it is not possible to conclude from the IRC study that from an asymmetric transition state the disrotatory stereochemistry is favored. To clarify this question, Mann and Hase carried out dynamical calculations [27, 28]. The study of dynamics was carried out with a limited number of trajectories, but the authors conclude that with a large ensemble of trajectories one may predict no stereochemical preference for the ring opening. To explain this result, the authors argue that the possible existence of a VRI point along the PES region is associated to the ring-opening process. More recently, an

[^1]Fig. 1 Orbitals of the minimum $\mathrm{M}_{c}: \sigma$ double-occupied, $n$ single-occupied, $\sigma^{*}$ un-occupied


Fig. 2 Orbitals of the minimum $\mathrm{M}_{a}: \pi_{1}$ double-occupied, $\pi_{2}$ single-occupied, $\pi_{3}$ un-occupied


IRC analysis of this reaction concludes that there exists a VRI point near the IRC curve, and this VRI point is located in a ridge that divides the initial valley in two valleys each one related to the allyl radical [29]. Although the result reveals the existence of the hypothetical VRI, it is not complete conclusive because the VRI point is located in a non-symmetric bifurcation valley due to the fact, as explained above, that the transition state is not symmetric. Recently, it has been proved computationally and experimentally for ion alkali halo-alkane molecular systems that the existence of VRI points and bifurcating valleys on their PES have strong relations to their dynamics behavior and product ratios of this type of molecular systems [30].

The purpose of the present study is both to give a more detailed structure and features of the UHF PES (unrestricted Hartree-Fock) related to the ring opening of cyclopropyl radical and the potentiality of the NTs as a tool to view and analyze a PES. More specifically and with a concrete problem, we start with the following question: the localization of the transition state and the RP on the PES imposes the condition that the path at some points loses the symmetry of allyl and cyclopropyl radicals; this implies that there exist bifurcation points where it happens. How the NT curves perform the localization of these special pathways through the bifurcation points represents a real challenge. Note that we take NTs in valley regions as RP models.

As a starting point of this task, we have proved many initial guesses of NTs computed at the UHF/6-31G(d) level of theory (for reasons of cost) where at least some of them should be related to the conrotatory and disrotatory interconversion modes. However, we did not find NT pathways with the features related to the pure conrotatory or disrotatory motion. These results are in accord with the findings of Ref. [25]. An RP does not exist in the ground electronic state joining the minimums associated with allyl and

cyclopropyl radicals, and preserving pure conrotatory or disrotatory motion. It is equivalent that the thermal conrotatory and disrotatory reactions are forbidden.

For the present purpose, we take the UHF/6-31G(d) level of theory as a way to generate the PES associated with the electronic ground state of the system. This wave function is very flexible and not very expensive from a computational point of view, also its solutions reasonably approach to the corresponding CASSCF wave function solutions (complete active space SCF) [31, 32]. The same type of wave functions was used in Ref. [25, 27-29] in the study on the cyclopropyl ring-opening system, see also Figs. 1 and 2. As explained very well in Ref. [25] it is interesting and important to analyze at different points of the PES the expectation value of the spin-squared operator $S^{2},\left\langle S^{2}\right\rangle$, of the UHF wave function. This magnitude in the present system indicates the spin contamination of states of higher multiplicity (i.e. quartet, sextet, ...) to the doublet state multiplicity. Normally, in the points of the PES where $\left\langle S^{2}\right\rangle$ differs from the value of the pure doublet wave function namely, 0.75, the UHF wave function shows a spin contaminant mainly due to the quartet state. If one uses the single-configuration approach, then the wave function of the quartet state can be written as a linear combination of three Slater determinants where each determinant is characterized by three molecular orbitals. These three molecular orbitals are related in some way with the $\sigma$ double-occupied, $n$ single-occupied, and $\sigma^{*}$ unoccupied of cyclopropyl molecular orbital system, Fig. 1, and with the $\pi_{1}$ double-occupied, $\pi_{2}$ single-occupied, and $\pi_{3}$ un-occupied orbitals of the allyl $\pi$-molecular orbital system, Fig. 2. We refer to these molecular orbitals as $\phi_{1}$, $\phi_{2}$, and $\phi_{3}$ in a generic form. The $\sigma$ and $\sigma^{*}$ molecular orbitals are related to the CC bond of the cyclopropyl radical to be broken in the evolution to the allyl radical, they are located in this bond, whereas the $n$ orbital is the
single-occupied molecular orbital located at the methine group,- C 1 H 4 , of cyclopropyl molecular system.

According to these results, the ring-opening system is minimally well described by three electrons and three orbitals, namely, $\phi_{1}, \phi_{2}$, and $\phi_{3}$ whereas the rest of electrons acts as a core potential. Due to this fact, the CASSCF wave function capable to describe this reaction adequately should be constructed using three electrons distributed within the above three molecular orbitals and the rest of electrons acting as spectators [25].

Nevertheless, in this extensive study of the PES we have found points and regions where the above electronic description is not enough, in addition to the three mentioned electrons, some core electrons are necessary to obtain a correct electronic description. These points have been associated with intersection points of the ground state with the PES of an excited electronic state; for this compare a work for cyclopropyl iodide [33]. The BornOppenheimer approximation assumes that the electrons relax instantaneously to their lowest energy distribution. For intersecting PESs, the situation becomes different. Because the electrons are strongly quantum mechanical, the transition cannot be a slowly "cooling" but has to be an ultrafast process that occurs at molecular geometries where the electronic states are isoenergetic [34, 35]. The geometries constitute a conical intersection and can be thought as the TS in the relaxation of the electronic excited molecule. However, they are not isolated points. Rather, they are collections of geometries that form a higher dimensional seam.

The molecular structure and the orbitals computed at UHF/6-31G(d) level of theory for the minimums of cyclopropyl and allyl radical are given in Figs. 1 and 2. The internal coordinates are reported in appendix in Tables A1 and A 6. The atoms numbering in an opened structure is given in Fig. 3. This and the z-matrix are used throughout the present study. The GAMESS-US suite of programs has been used for the calculations [19, 20].

The used z-matrix is

| Cyclopropyl radical $<-->$ allyl radical |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| C1 |  |  |  |  |  |  |
| c1 |  |  |  |  |  |  |
| c2 | 1 | r1 |  |  |  |  |
| c3 | 1 | r2 | 2 | c2c1c3 |  |  |
| h4 | 1 | hc1 | 2 | hcc1 | 3 | dih1 |
| h5 | 2 | hc2 | 1 | hcc2 | 3 | dih2 |
| h6 | 3 | hc3 | 1 | hcc3 | 2 | dih3 |
| h7 | 2 | hc4 | 1 | hcc4 | 3 | dih4 |
| h8 | 3 | hc5 | 1 | hcc5 | 2 | dih5 |



Fig. 3 Ring opening of the cyclopropyl radical: numbering of atoms in the z-matrix

## 3 The cyclopropyl radical bowl

In Table 1, we report a condensed version of the coordinates of all interesting structures, beginning with the cyclopropyl radical minimum, $\mathrm{M}_{c}$, at UHF/6-31G(d) level of theory. The index c is used for points of the cyclopropyl part. This geometry structure belongs to the $\mathrm{C}_{s}$ point group of symmetry, and due to this fact the distance r 1 is equal to r2 and the dihedral angles satisfy the relation $\operatorname{dih} 2=-\operatorname{dih} 5$ and dih3 $=-$ dih4. With the dihedral that involves the H 4 atom, this set of coordinates is the most important and relevant in the changes of the molecule. The remainder of pairs of geometry parameters also preserves the $\mathrm{C}_{s}$ symmetry relation. The numbering of the two methylenes is interchangeable. Thus, in the present study we will very often treat the corresponding pairs of coordinates, also if the molecular structure does not belong to the $\mathrm{C}_{s}$ point group of symmetry, but they hold their connection to the other corresponding part.

From the minimum $\mathrm{M}_{c}$, two "reaction directions" emerge. One is the 1 D transit of the single methine hydrogen, H 4 , through the symmetry plane of the carbons, thus the dihedral angle goes through $\pm 180^{\circ}[36,37]$. There is an $\mathrm{SP}_{H 4}$ of the PES, see in the appendix Table A2. The remainder of the coordinates stays nearly unchanged at the minimum values.

The other possibility is the ring opening of the $\mathrm{M}_{c}$ structure by increasing the angle C 2 C 1 C 3 , which is equivalent to the bond breaking C2C3. The bond angle C 2 C 1 C 3 will be our reference coordinate in the figures, so to say, it will be our reaction coordinate. The reaction may start in the symmetric subspace of all corresponding coordinate pairs of minimum $\mathrm{M}_{c}$, thus there is, at least a 3D, subspace of important possibilities for the pairs $\mathrm{r} 1=\mathrm{r} 2, \operatorname{dih} 2=-\operatorname{dih} 5$, and $\operatorname{dih} 3=-\operatorname{dih} 4$. We note that the z-matrix has 18 internal coordinates that we can decompose as follows, seven pairs each pair being related by the $\mathrm{C}_{s}$ symmetry, the C 2 C 1 C 3 bond angle and the three coordinates of the H 4 atom. Note: if the dihedrals

Table 1 Energies and special coordinates of stationary and VRI points on the PES of cyclopropyl radical opening to allyl radical (for UHF 6-31G(d))

| point | Energy au | $\begin{aligned} & \mathrm{r} 1^{\mathrm{a}} \\ & \mathrm{C} 2 \mathrm{C} 1 \end{aligned}$ | r2 <br> C1C3 | $\begin{aligned} & \angle \\ & \mathrm{CCC} \end{aligned}$ | dih1$\mathrm{H} 4^{\mathrm{b}}$ | dih2/dih4 |  | dih3/dih5 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  | H5/H7 ${ }^{\text {b }}$ |  | H6/H8 |  |
| $\mathrm{M}_{c}$ | -116.41492 | 1.471 | 1.471 | 63.5 | $\pm 120 .{ }^{\text {c }}$ | 106.1 | $-107.5$ | 107.5 | -106.1 |
| $\mathrm{SP}_{H 4}$ | -116.40935 | 1.456 | 1.456 | 63.4 | $\pm 180.0$ | 106.3 | -106.3 | 106.3 | -106.3 |
| $\mathrm{VRI}_{H 4}$ | -116.36293 | 1.464 | 1.464 | 84.7 | -148.0 | 97.0 | -99.3 | 99.3 | -97.0 |
| $\mathrm{SP}_{i 2}$ | -116.34794 | 1.448 | 1.451 | 90.8 | -179.9 | 89.1 | -89.2 | 89.2 | -89.2 |
| $\mathrm{VRI}_{c}$ | -116.37773 | 1.451 | 1.451 | 79.9 | -132.9 | 108.5 | $-100.0$ | 100.0 | -108.5 |
| $\mathrm{SP}_{c a}$ | -116.37268 | 1.416 | 1.471 | 86.7 | -127.6 | 114.8 | -71.5 | 87.0 | -95.7 |
| $\mathrm{SP}_{c a}$ |  | 1.471 | 1.416 | 86.7 | -127.6 | 95.7 | -87.0 | 71.5 | -114.8 |
| $\mathrm{SP}_{c a}$ |  | 1.471 | 1.416 | 86.7 | $127.6{ }^{\text {d }}$ | 87.0 | -95.7 | 114.8 | -71.5 |
| $\mathrm{SP}_{c a}$ |  | 1.416 | 1.471 | 86.7 | $127.6{ }^{\text {d }}$ | 71.5 | -114.8 | 95.6 | -87.0 |
| $\mathrm{VRI}_{c a}$ | $-116.37565$ | 1.401 | 1.488 | 93.6 | -127.2 | 120.7 | -58.4 | 76.3 | -89.2 |
| $\mathrm{VRI}_{c a}$ |  | 1.488 | 1.401 | 93.6 | -127.2 | 89.2 | -76.3 | 58.4 | -120.7 |
| $\mathrm{VRI}_{r}$ | -116.37204 | 1.403 | 1.503 | 97.8 | -140.9 | 97.8 | -76.1 | 80.6 | -82.6 |
| $\mathrm{VRI}_{r}$ |  | 1.503 | 1.403 | 97.8 | -140.9 | 82.6 | -80.6 | 76.1 | -97.8 |
| $\mathrm{VRI}_{s}$ | $-116.30119$ | 1.458 | 1.458 | 108.7 | -180.0 | 82.3 | -82.3 | 82.3 | -82.3 |
| $\mathrm{SP}_{d 1}$ | -116.37394 | 1.478 | 1.470 | 117.0 | -180.0 | 78.9 | -78.9 | 82.2 | -82.2 |
| $\mathrm{SP}_{d 1}$ |  | 1.470 | 1.478 | 117.0 | -180.0 | 82.2 | -82.2 | 78.9 | -78.9 |
| $\mathrm{VRI}_{r l}$ | $-116.42231$ | 1.316 | 1.470 | 115.5 | -172.0 | 168.6 | -7.8 | 102.8 | -85.1 |
| $\mathrm{VRI}_{i 2}$ | $-116.37769$ | 1.459 | 1.456 | 124.1 | -176.2 | 93.8 | -85.3 | 85.5 | -93.7 |
| $\mathrm{SP}_{i 2}$ | $-116.37782$ | 1.473 | 1.473 | 124.7 | -180.0 | 91.6 | -91.6 | 91.6 | -91.6 |
| $\mathrm{VRI}_{a}$ | -116.44217 | 1.407 | 1.407 | 128.3 | -179.9 | 152.2 | -39.1 | 39.1 | -152.2 |
| $\mathrm{SP}_{a a}$ | $-116.43851$ | 1.326 | 1.479 | 124.8 | -180.0 | 180.0 | 0.0 | 97.3 | -97.3 |
| $\mathrm{SP}_{a a}$ |  | 1.479 | 1.326 | 124.8 | -180.0 | 97.3 | -97.3 | 180.0 | 0.0 |
| $\mathrm{M}_{a}$ | -116.46810 | 1.390 | 1.390 | 124.6 | -180.0 | $180.0{ }^{\text {e }}$ | 0.0 | 0.0 | -180.0 |
| $\mathrm{M}_{a}$ |  | 1.390 | 1.390 | 124.6 | -180.0 | 0.0 | -180.0 | 180.0 | 0.0 |

${ }^{\text {a }}$ See Fig. 3 for numbering of the atoms. A subscript $c$ is used for the symbols of the points in the cyclopropyl radical bowl, but a subscript $a$ is used for points in the allyl radical bowl; see the text for further symbols
${ }^{\text {b }} \mathrm{C} 1 \mathrm{H} 4$ is the methine, C 2 H 5 H 7 is methylene 1 , and C 3 H 6 H 8 is methylene 2
c Transition of sign changes the H atom numbering only
${ }^{d}$ For comparison, the two other symmetrical dih1-forms are added. Corresponding symmetrical forms also concern other points
${ }^{e}$ The first $M_{a}$ is the disrotatory minimum, the second one the conrotatory
are involved, then a change of the molecule in the symmetry subspace $\mathrm{C}_{s}$ is a disrotatory change of the two methylenes.

We try to find the true singular NT in the 3D subspace in order to explore the PES from $\mathrm{M}_{c}$ to the bifurcation point. To do this, we start with some test calculations. One should take into account that NTs can turn back to the minimum, after passing a so- called turning point (TP) [38]. Or, if they start at the SP, they can turn up and go lost anywhere in the mountains of the PES. Due to this fact, one needs at least an NT which connects $\mathrm{M}_{c}$ with the $\mathrm{SP}_{c a}$ (their coordinates are in appendix Table A1). The true iteration of $\mathrm{VRI}_{c}$ by our program can begin if an NT with a kink indicates a neighborhood of the bifurcation. The present case shows a special difficulty, because the criterion $|\mathbf{A g}|$ does not show
a tendency to go to zero [39] also in a small neighborhood of the $\mathrm{VRI}_{c}$. At least, it emerges that the function $|\mathbf{A g}|$ is very steep along the found singular NT to $\mathrm{VRI}_{c}$, which is depicted in Fig. 4.

At this point, we make the following note with respect to the representation. We use both dih2 and - dih5 up to $\mathrm{VRI}_{c}$ because they are identical. However, after the VRI point, they bifurcate and go through different values and describe the behavior of the unsymmetric reaction path. (If we accept that the singular NT can be a description of an RP.) But we could renumber the molecule; then dih2 becomes -dih5, and vice versa. An analog argument concerns the coordinates dih3 and -dih4, as well as the two distances of the C -atoms, r 1 and r 2 . The representation of Fig. 4 (left panel) means that we have the two


Fig. 4 The VRI point, $\mathrm{VRI}_{c}$, of the pre-transition-state symmetry break for the ring opening of cyclopropyl radical indicated by the singular NT and its bifurcation point, see text. The cross of grid lines is the exact location of $\mathrm{VRI}_{c}$. Left the representation uses the doubling of all points for coordinates dih2, as well as -dih5, to the same bond

angle coordinate C 2 C 1 C 3 . The dihedrals are interchangeable. The associated H atoms are one possibility. Symbol asterisk is a point at the end of one PES, and an electronic intersection happens with another PES. Right The same singular NT through $\mathrm{VRI}_{c}$ for coordinates dih3 and -dih4. The inlay is the geometry of $\mathrm{VRI}_{c}$
initial molecular symmetry of $\mathrm{M}_{c}$ to arrive at the corresponding $\mathrm{SP}_{c a}$ [25].

When the upper branch of the left panel of Fig. 4 concerns coordinate dih2 being related to H5, and consequently, when the lower branch of the same panel in Fig. 4 concerns with the coordinate - dih5 which belongs to H8, and dih2 becomes larger but dih5 becomes smaller, then the evolution from $\mathrm{VRI}_{c}$ to $\mathrm{SP}_{c a}$ occurs through a conrotatory change. The two hydrogens move simultaneously to the left-hand side in Fig. 1. With a similar observation and reasoning, when the upper branch of the right panel of Fig. 4 concerns with the coordinate dih3 being related to H6, and when the lower branch of the right panel of Fig. 4 is coordinate - dih4 which belongs to H 7 , and both dih3 and dih4 become smaller, then the evolution from $\mathrm{VRI}_{c}$ to $\mathrm{SP}_{c a}$ occurs through a (something disturbed) disrotatory character. The two hydrogens turn to an outer space, in Fig. 1. The view of the whole process means that the methylene 1 related to C 2 rotates clockwise, whereas the methylene 2 related to C 3 is compressed, because H 6 and H 8 approach in a concerted kind. Note: the "four" SPs in Fig. 4 concern one and the same structure of the cyclopropyl radical, correspondingly, to the four dihedral coordinates used.

It is important at this point to note that here emerges the intersection problem of excited PESs within the UHF methodology for the first time: An intersection point along the singular NT is indicated by a * symbol in the left panel of Fig. 4. It says that the structure, number, and occupancy of UHF natural orbitals with fractional occupancy has changed. Generally, we use the UHF methodology rather than CASSCF; however, the UHF methodology is not appropriate in the * region. We cut the calculation of the NT at this point. An expected symmetric SP of index 2 could not be reached going along this branch, because of the intersection with another PES.


Fig. 5 Left The VRI point, $V_{R I}$ marking the border between two regions of attraction of the interesting SPs of cyclopropyl radical. It is indicated by the singular NT and its bifurcation point. The grid lines cross depicts to the location of $\mathrm{VRI}_{H 4}$. Symbols asterisk are points at

the end of this UHF PES. For comparison, the singular NT of $\mathrm{VRI}_{c}$ is also given by thin points, see Fig. 4. The inlay is the geometry of $\mathrm{VRI}_{H 4}$. Right $\mathrm{VRI}_{H 4}$ for coordinate dih1. The additional left NT is a regular one between $\mathrm{M}_{c}$ being left top, and $\mathrm{SP}_{H 4}$ being left bottom


Fig. 6 Left a regular NT to $\mathrm{SP}_{i 2}$ for coordinates dih2 and - dih5. Compare the caption of previous Figures. The $\mathrm{VRI}_{c}$ system is given in the left panel by thin points, for comparison. Right The same regular NT to $\mathrm{SP}_{i 2}$ for coordinate dih1

The border between the regions where NTs going to $\mathrm{SP}_{H 4}$, or others going to an $\mathrm{SP}_{c a}$, is marked again by a singular NT which bifurcates at a symmetric bifurcation point in $\mathrm{C}_{s}$, the $\mathrm{VRI}_{H 4}$, see appendix Tables A2. It is demonstrated in Fig. 5 where additionally to the dihedrals dih2 and -dih5, also the behavior of dih1 is shown. It is clear that here the bifurcation mainly concerns the coordinate dih1 in a plane with the bond angle C 2 C 1 C 3 . There is experimental evidence that the lone H 4 of the C 1 H 4 bond can tunnel through the barrier of $\mathrm{SP}_{H 4}$, if the cyclopropyl radical is vibrationally exited [37]. One should hold in mind the fact along the full ring opening.

Looking at the left panel of Fig. 5, the singular NT through $\mathrm{VRI}_{H 4}$ meets "with an accident" because the two bifurcating branches are no mirror pictures, like in other bifurcations. The character of the $\mathrm{VRI}_{H 4}$ is of the mixed type, thus it is neither vpVRI nor rpVRI. The two upper branches in Fig. 5 are valleys of the PES, but the two lower branches are ridges of index one. The two kinds of curves have a kink and "touch with the kink" at the VRI point [15, 17]. However, the name VRI may be correct: if we look at
the right panel in Fig. 5, there the straight line from $\mathrm{M}_{c}$ in the left upper corner to the right lower region describes a valley-ridge inflection, as well as the two branches which connect the $\mathrm{SP}_{H 4}$ in the left lower corner with the right upper region describe a ridge-valley transition.

The two right branches meet after the $\mathrm{VRI}_{H 4}$ at points that intersect with the PES of an excited electronic state. That is the reason that a possible connection to $\mathrm{SP}_{c a}$ could not be shown.

In Fig. 5 we also show the location of an SP of index two, $\mathrm{SP}_{i 2}$ on the PES of the $\mathrm{M}_{c}$ bowl. The Hessian matrix of this SP was computed using GAMESS-US [19, 20], and the corresponding diagonal representation shows two negative eigenvalues. Figure 6 shows a regular NT connecting the $\mathrm{SP}_{c a}$ with this $\mathrm{SP}_{i 2}$ of index two, for the internal coordinates, compare Table A2 in the appendix. One may imagine the start point after the $\mathrm{SP}_{c a}$ at the allyl side, the NT then crosses the SP as a minimum on the way, and goes uphill the energy surface, see also Fig. 17 below. In the left panel of Fig. 6 we see that the NT has a returning behavior for the angle C 2 C 1 C 3 , at $\approx 81.5^{\circ}$. Note that the $\mathrm{SP}_{i 2}$ is in


Scheme 1 Connection scheme of the cyclopropyl radical bowl. $\mathrm{M}_{c}$ : cyclopropyl radical minimum. $\mathrm{SP}_{c a}$ : TS of the ring opening of cyclopropyl radical to allyl radical. $\mathrm{SP}_{i 2}$ : SP of index 2 where $\mathrm{H}_{4}$ is in the CCC plane. The two boxes concern the same point. VRI valleyridge inflection point of a singular NT between $\mathrm{M}_{c}$ and corresponding SPs. Asterisk depicts an intersection with another electronic surface; a systematic search by NTs breaks down. Corresponding NT arrows are hypothetical. The lower line of VRI points is outside the cyclopropyl radical bowl
the region of symmetry coordinates $\mathrm{C}_{s}$, like $\mathrm{M}_{c}$; however, it is not totally symmetric like $\mathrm{M}_{a}$ (the index a is used for the allyl radical minimum) with $\mathrm{C}_{2 v}$ symmetry. For comparison, we also include the singular NT of $\mathrm{VRI}_{c}$ in Fig. 6. Its branches through the $\mathrm{SP}_{c a}$ should go, more or less, along the valley floor of the SP.

We also computed an NT which connects the $\mathrm{SP}_{i 2}$ with the $\mathrm{SP}_{H 4}$, which is not shown in the Figs. 5 and 6. In Scheme 1 we collect the complete set of treated relations in the cyclopropyl radical bowl, compare also Fig. 17 below.

## 4 The $\mathbf{S P}_{c a}$ and the $\mathrm{VRI}_{c a}$ in the $\mathbf{S P}$ valley

The $\mathrm{SP}_{c a}$ associated with the ring opening of the cyclopropyl radical into the allyl radical is well known [25]. The coordinates at the UHF/6-31G(d) level of theory used in the present study are reported in appendix Table A1. The SP is not symmetric like the cyclopropyl or allyl radicals, compare Fig. 4. After the SP, with a further ring opening, there should be the post-TS bifurcation because at a possible end of the pathway one expects to find one of the diverse versions of an $\mathrm{SP}_{a a}$ in the allyl radical bowl, see
below; and due to this fact the downhill pathway flowing into the $\mathrm{SP}_{a a}$ region from $\mathrm{SP}_{c a}$ may be a ridge. At the beginning, the pathway is a valley. As an immediate consequence a VRI point must exist (Fig. 7).

To find this $\mathrm{VRI}_{c a}$ point was especially complicate, because the corrector of our method to follow NTs [8] does not properly work in the whole region around this VRI point. (Maybe because the determinant of the Hessian is very small, or because of any other reason.) See for comparison an example reported in [18] for the case of HCN with an analog but still worse behavior. Thus, we have to calculate the given NTs by predictor steps only using Branin-Euler steps along an equation of the tangent $\mathbf{x}^{\prime}= \pm$ Ag. However, this does not mean to be a breakdown. Because near the VRI point, regular NTs are usually strong curvilinear but the searched VRI is usually outside of the convex region. Tangent steps in direction $\mathbf{x}^{\prime}$ along such a curvilinear NT make their approximation error into the "right" direction. They lead nearer to the VRI point. If the VRI point is reached, the emergence of a cross of branches of the singular NT demonstrates the final success, see Figs. 8 and 9.

The electronic structure is very interesting because it represents the breakdown of the electronic state of the cyclopropyl structure, and it is the beginning of the unsymmetric region. The three important natural orbitals are shown in Fig. 7.

Every VRI system of a singular NT consists of four branches. The character of the single branches determines the character of the VRI point. The character of the $\mathrm{VRI}_{c a}$ is a branching of the uphill leading ridge from an hypothetical $\mathrm{SP}_{a a}$ region (being at the right-hand side, out of the Fig. 8) into two ridges going strong further uphill, and the one valley in between, which leads to the $\mathrm{SP}_{c a}$. The character is an rpVRI of the chemically "interesting" case: between the bifurcating ridge-branches is the valley. (See below for singular NTs in not so "interesting" cases.) Thus, the ring-opening valley from $\mathrm{SP}_{c a}$ downhill ends at the VRI point. There are no bifurcating valleys downhill. Some authors name the situation of such an ending valley a "dangerous" bifurcation [41], or a bifurcation with indeterminate outcome, because further downhill after the VRI point, there the valley ground line disappears. After the $\mathrm{VRI}_{c a}$, only the ridge-pathway continues further downhill. And going downhill on the PES, the ridge here shows a

Fig. 7 Natural orbitals of the VRI $_{c a}$ : double-occupied with 1.723 occupation, singleoccupied with 1.0 occupation, and un-occupied with 0.277 occupation



Fig. 8 VRI $_{c a}$ point behind the transition state, $\mathrm{SP}_{c a}$, of the ring opening of the cyclopropyl radical. PES by UHF/6-31G(d). Left pieces of regular NTs around, as well as of the singular NT through the $\mathrm{VRI}_{c a}$. Coordinates: angle between the C -atoms, and the two replaceable dihedrals dih2 and -dih5. The cross of the three grid

lines is the location of the VRI point. The asterisk symbol again indicates the electronic intersection of a second PES of the allyl radical side of the problem. There, the current NT ends. The indicated here continuation is an artifact, at all. Right coordinates are the angle between the C -atoms, and the two dihedrals $\operatorname{dih} 3$ and -dih4, see text

Fig. 9 VRI $_{c a}$ for Left bond length coordinates r1 and r2 and Right dihedral dih1, see caption of Fig. 8. The inlay is $\mathrm{VRI}_{c a}$

new effect, a "double bend" which indicates a road irregularity of the PES. A proposed explanation is added in the next Section.

The evolution pathway from $\mathrm{SP}_{c a}$ to $\mathrm{VRI}_{c a}$ includes quasi no change in the distance of the dihedral angles dhi3 and -dih4 in the right panel of Fig. 8. The two NTs are parallel. However, there is a small increase in the difference of dih2 and -dih5 (left panel of Fig. 8). The upper dihedral is increasing and the lower is decreasing. It means, the two upper hydrogens, H6 and H7 in Fig. 3, move disrotatory, where the two lower hydrogens, H8 and H5 in Fig. 3, move conrotatory along the left branch of the singular NT. The situation is changing drastically after the VRI point.

There, after $\mathrm{VRI}_{c a}$ at the right-hand side in every of the two panels of Fig. 8, the branch of the singular NT is the crest of a ridge downhill. The double bend of this branch is quite different for either the methylene of C 2 (dih2, -dih4), or the methylene of C3 (-dih5, dih3). Near $98^{\circ}$ of the angle C 2 C 1 C 3 , there is a returning moment: methylene of C 2 only stops, goes back, stops again and goes further to its aim. However, methylene of C3 fulfills a turn into a conrotatory direction. The distance (dih2 to -dih5) changes little; however, the distance (dih3 to -dih4) modifies very much.

Note that after the double bend of the NT there is the electronic intersection between PESs. One may speculate that this computed UHF PES will get ready for the crash by the tremor of the methylenes.

## 5 The effect of bumps in the road

There is a sudden change between disrotatory and conrotatory motion along some branches of the singular NT in a larger region of the PES, see Fig. 8 for the ridge after the $\mathrm{VRI}_{c a}$ point. We demonstrate by a simple two-dimensional test surface, how such a road irregularity of NTs may emerge on a molecular PES. The left panel of Fig. 10 is calculated by the Mathematica ${ }^{\circledR}$ program. Such road irregularities emerge for NTs on the upper PES after the $\mathrm{SP}_{c a}$ and can be seen when the curves are represented by bond angle C2C1C3 and dihedral dih2 shown in the right panel of Fig. 10. The possible additional small ridge on the PES then falls directly before the intersection of the PES related to another electronic state. The symbol * indicates here a precipice to a lower PES in the energy. We found that the UHF method is possibly not fully appropriate for the regions where the cliffs located between the plateau region of cyclopropyl radical

Fig. 10 Left 2D-test surface with an SP at $(0,0)$, the SD from there (fat points), and a road irregularity for 5 NTs (dashes) passing the SP. The double bend of two connected turns that travel in alternate directions emerges by a very small, additional nose on the PES. It is produced by the function $\operatorname{Exp}[-20(x+0.5 y$ $\left.-1)^{2}\right]$ which is added to the saddle surface $3.5\left(x^{2}-y^{2}\right)$. The steepest descent is much less disturbed than the NTs. Right Some NTs after $\mathrm{SP}_{c a}$ with road irregularities


side, and the ground region of the allyl radical appear. Surprisingly, UHF does not converge in a suitable solution there. More specifically, the UHF solution gives a broken solution ( $\left\langle S^{2}\right\rangle$ is different from 0.75); however, it is not the correct solution that still is or should be a broken solution in the form that it is higher to 0.75 . The UHF wave function possesses a set of UHF solutions in the region of the pronounced precipice, and one should take the correct solution: it is a task which is not easy. In fact the algorithm implemented in GAMESS-US [19, 20] package program to find the UHF in this case always converges to the non-appropriated solution. More explicitly, many times it converges either to the expected value for a pure doublet wave function solution ( $\left\langle S^{2}\right\rangle=0.75$ or near it) or to the highest broken solution.

The intersection of the PESs itself is of the type "product side" [42], or, with another classification, of the type "peaked intersection" [35].

For some NTs, the nodes after the precipice indicate a continuation of the NT: that can be done by using the current gradient for the search direction of the new NT. Of course, it is not the continuation of the former NT, but it is a new one starting at the same point in coordinate space but with electronic structure of the lower surface. The gradients are quite different, usually, at the intersection of the two different surfaces. A systematic use of the NT theory is not possible if one overcomes the precipice.

## 6 Singular ridges on the upper $\operatorname{PES}\left(\mathrm{VRI}_{r}, \mathrm{VRI}_{s}\right)$

In the right panel of Fig. 10, the next VRI point still near the $\mathrm{SP}_{c a}$ is shown named $\mathrm{VRI}_{r}$. It is on the upper UHF PES,
the PES of the cyclopropyl side. The UHF energy is -116.37204 au , see also Fig. 17 below, with $\left\langle\mathrm{S}^{2}\right\rangle=1.517$. The occupancies of the natural orbitals $\phi_{1}$ (being the orbital 11 in the natural orbital list), $\phi_{2}$ (being the orbital 12 ), and $\phi_{3}$ (being the orbital 13) are $1.507,1.0$, and 0.493 , correspondingly. The occupancy reflects a strong delocalization in the electronic structure. It is due to the bond breaking between the atoms C2 and C3. A control calculation by the CASSCF method already shows a slow transition to a mix between $\pi$ - and $\sigma$-electronic structure. The CASSCF energy of this point is -116.38154 au. The occupancies of the most important natural orbitals are for $\phi_{1}$ (being the orbital 11 in the natural orbital list), $\phi_{2}$ (being the orbital 12), and $\phi_{3}$ (being the orbital 13) 1.1292, 1.0 , and 0.8708 , correspondingly. The $S^{2}$ operator commutes with the CASSCF Hamiltonian operator. It implies that for this energy functional the expectation value $\left\langle S^{2}\right\rangle$ is always 0.75 for a doublet state.

The CASSCF result for the second root (electronic state) not optimized is located at $\approx 30 \mathrm{kcal} / \mathrm{mol}$ with respect to the ground state. It supports that the "upper" UHF wave function is appropriate in this region. But the $\mathrm{VRI}_{r}$ point has more broken bonds than $\mathrm{SP}_{c a}$, compare [25]. The structure of the orbitals can be found in Fig. 11. These orbitals are a mix between $\sigma$ - and $\pi$-electronic structure. The movement of the first two eigenvectors of the Hessian matrix corresponds to a broken "symmetry C2C1C3 plane" in both directions, see below Fig. 22.

The $\mathrm{VRI}_{r}$ defines a ridge system that divides NTs going downhill to $\mathrm{SP}_{a a}$ (which is not shown) and other NTs going uphill into non-chemical mountains of the upper PES. We named it $\mathrm{VRI}_{r}$, see Table A3 in the appendix, because it is

Fig. 11 Structure of CASSCF natural orbitals $\phi_{1}$ to $\phi_{3}$ (see text) at $\mathrm{VRI}_{r}$



Fig. 12 Left branches of the singular NT through $\mathrm{VRI}_{r}$ behind the transition state, $\mathrm{SP}_{c a}$, on the upper electronic PES. Coordinates are the angle between the C -atoms, and the two dihedrals dih2 and -dih5. The cross of the three grid lines is the location of the VRI point. The

asterisk symbol indicates again an electronic intersection of a second PES of the allyl radical side of the problem. The inlay is the shape of VRI $_{r}$. Right coordinates are the angle between the C-atoms, and the two dihedrals dih3 and -dih4. See text
behind the $\mathrm{SP}_{d 1}$ meet again at an intersection with another electronic PES indicated by the * symbol.

A next VRI point on the upper surface is one between the $\mathrm{SP}_{H 4}$ of the $\mathrm{M}_{c}$ bowl and again the region of the $\mathrm{SP}_{d 1}$. The point is named $\mathrm{VRI}_{s}$ because it is in the symmetric subspace $\mathrm{C}_{s}$ like the $\mathrm{VRI}_{c}$. Its singular NT is a quasistraight line in the coordinate space, and the angle dih1 remains throughout near $-180^{\circ}$. It ends near the $\mathrm{SP}_{d 1}$ in an electronic intersection. The internal coordinates of $\mathrm{VRI}_{s}$ are given in Table A3 in the appendix. The fractional occupancies of the natural orbitals are $2.0,1.0$, and 0.0 . Thus, the avoided crossing seems to be far away. The UHF energy of $\mathrm{VRI}_{s}$ is very high: -116.30119 au, see also Fig. 17 below. We think that it has no greater chemical meaning, without that to mark a high symmetry border between the chemical more interesting, unsymmetric regions of the PES. The singular NT is a ridge ascenting from $\mathrm{SP}_{H 4}$. It bifurcates at $\mathrm{VRI}_{s}$ into a symmetric ridge of index two, and two unsymmetric side branches (of index 1) which also quickly end at an intersection seam. This ridge ascenting from $\mathrm{SP}_{H 4}$ is another one than the ridge ascending from $\mathrm{SP}_{H 4}$ to $\mathrm{VRI}_{H 4}$, compare Fig. 5.

## 7 Singular ridges on the lower $\operatorname{PES}\left(\mathrm{VRI}_{r l}\right.$, VRI $\left._{i 2}\right)$

One could speculate that a ridge on the lower PES, rising from $\mathrm{SP}_{a a}$ in direction of $\mathrm{SP}_{c a}$ under a decreasing bond angle C2C1C3, has to meet a VRI point where it transforms into a valley. But we could not find such a corresponding singular NT. Contrary, we found an uphill ridge from $\mathrm{SP}_{a a}$ which meets a $\mathrm{VRI}_{r}$ of the character rpVRI, but of a higher index, see Fig. 13. We named it $\mathrm{VRI}_{r l 1}$ with " r " because it is a bifurcation of ridges, " 1 " for the lower PES, and " 1 " for the first one, see below. Its coordinates are in Table A3 in the appendix. The character is a bifurcation of a usual


Fig. 13 Left singular NTs through two $\mathrm{VRI}_{r l}$ one after the other, on the lower PES, for coordinates dih2 and -dih5 against the angle C2C1C3. Right the same NTs for coordinates dih3 and -dih4. At the left-hand side of every panel are branches of the singular NT of the VRI ${ }_{c a}$

Fig. 14 Natural orbitals of the $\mathrm{SP}_{i 2}$

ridge of index one into two ridges of index one, and a ridge of index two in between, so to say a two-dimensional ridge on the 18 -dimensional PES. Later, a further bifurcation happens at a $\mathrm{VRI}_{r l 2}$ of the central ridge of index two into two ridges of index two, and a central ridge of index three, thus a ridge in tree dimensions. We guess that it is of less interest for chemical interpretations. Because, also its energy is high in the PES mountains. If one compares the corresponding branches with the system of branches of the $\mathrm{VRI}_{c a}$, shown on the left-hand side by thinner points, it is clear that the two systems of VRI branches do not really fit together, though some single branches seem to meet.

On the central ridge of $\mathrm{VRI}_{r l 1}$, the dihedrals dih2 and -dih5 hold their distance over a ring closure of $20^{\circ}$ of the bond angle C2C1C3 where they slightly decrease. The dihedral dih3 is constant, and dihedral -dih4 increases a little. The singular NT of the $\mathrm{VRI}_{r l 1}$ corresponds to a quasidisrotatory ring closure if it starts at $\mathrm{SP}_{a a}$.

The full symmetric $\mathrm{SP}_{i 2}$ of index 2 is also on the lower PES, however, already of an high energy of -116.37782 au . Its character is of $\sigma$-type radical electronic state $\left({ }^{2} \mathrm{~A}^{\prime}\right)$. It is in the coordinate space near to the $\mathrm{SP}_{d 1}$ being on the upper PES, compare Fig. 12. Its character is of $\pi$-type radical electronic state (almost a ${ }^{2} \mathrm{~A}^{\prime \prime}$ because this SP is $\mathrm{C}_{1}$ symmetry). The UHF coordinates of $\mathrm{SP}_{d 1}$ and $\mathrm{SP}_{i 2 a}$ are in Tables A4 and A5 in the appendix. Possible NTs between the two SPs cannot seriously be followed over the electronic intersection.

The SP structure of index 2 may clarify many aspects of the PES at first sight. The first eigenvector of the Hessian matrix at the point with the negative eigenvalue

corresponds to the conrotatory motion while the second one is disrotatory. We have done a comparison of the UHF geometry and their Natural Orbitals by a single-point calculation using the $\operatorname{CAS}(3 \times 3)$ wave function, and it shows two almost degenerate electronic solutions. Maybe we are here in front of an intersection between the ground and the excited radical electronic states. It is important to take into account the conrotatory and disrotatory motion taking place in the excited state, see [25], and see the Section to the steepest descent from $\mathrm{SP}_{d 1}$ below. The geometry of this SP of index 2 belongs to the $\mathrm{C}_{2 v}$ point group of symmetry, see Fig. 14, the picture of the CAS Natural Orbitals with occupancies $0.774,1.0$, and 1.227, correspondingly.

Starting an NT on the lower PES at the $\mathrm{SP}_{i 2}$, this NT carries the index 2 from the beginning. We have also done that computational experiment hoping to loose some index. However, also such a singular NT, see Table A5 in the appendix, and Fig. 15, meets a $\mathrm{VRI}_{i 2}$ with an rpVRI character of higher index. Here, the branch between $\mathrm{SP}_{i 2}$ and the $\mathrm{VRI}_{i 2}$ is a ridge of index 2 . It bifurcates in two ridges of index 2 , and a ridge of index 3 in between. It may be of lower chemical interest. The branch between $\mathrm{SP}_{i 2}$ and the $\mathrm{VRI}_{i 2}$ is still symmetric for the coordinate pairs dih2/ -dih5 and dih3/-dih4, correspondingly. The three uphill branches after the VRI point then show a splitting of these pairs of coordinates. Again, the symmetry of $\mathrm{SP}_{i 2}$ is broken after the $\mathrm{VRI}_{i 2}$. The three forks of the ridge-pitchfork then rise up to non-interesting mountain regions of the PES. We guess that many further VRI points of higher index scrimmage in the region between $\mathrm{SP}_{a a}$ and $\mathrm{SP}_{c a}$.


Fig. 15 Left singular NT through $\mathrm{VRI}_{i 2}$, for coordinates dih2 and -dih5 against the angle C2C1C3. Right the same NT for coordinates dih3 and -dih4. Thin points are from the right branch of a former


Fig. 16 Left singular NT through $\mathrm{VRI}_{a}$ for coordinates dih2 and -dih5 against the angle C 2 C 1 C 3 . The small distances between $\mathrm{SP}_{i 2}$ and $\mathrm{SP}_{a a}$, as well as between $\mathrm{M}_{a}$ and $\mathrm{SP}_{a a}$, come from the projection only which is used in the figure. The upper inlay is $\mathrm{SP}_{i 2}$, the lower inlay is $\mathrm{VRI}_{a}$. Right the same NT for a direct comparison of the

## 8 The allyl radical bowl

The allyl radical minimum, $\mathrm{M}_{a}$, is a well-known structure at many levels of theory [25], see also appendix, Table A6 for the full coordinates, and Table 1 for the reduced version.

Interesting points in the bowl of the minimum often belong to totally symmetric structures, see Fig. 2. A singular NT which starts at $\mathrm{M}_{a}$ in the totally symmetric subspace of symmetry coordinates meets a $\mathrm{VRI}_{a}$ point where the valley bifurcates into two valleys to two SPs named $\mathrm{SP}_{a a}$, and a ridge in between to an SP of index two, $\mathrm{SP}_{i 2}$, see Fig. 16. The character of the VRI is a vpVRI. The UHF and CAS( $3 \times 3$ ) calculations at $\mathrm{VRI}_{a}$ coincide very much. The so-called "active space" (the electrons and orbitals which are strongly correlated) are those of the allyl radical minimum. The analysis of the eigenpair with null eigenvalue and orthogonal to the gradient vector reveals that it wants to move the

system of the singular NT of the $\mathrm{VRI}_{r l}$, for comparison. Note that $\mathrm{SP}_{i 2}$ and $\mathrm{SP}_{a a}$ are so closely related only in this projection

coordinates dih2 and dih5. It shows the symmetry relation of the branches of the NT. (It may also give an impression of the importance of an appropriate projection for the understanding of a bifurcation.) The inlays are versions of $\mathrm{SP}_{a a}$
structure toward a methylene orthogonal to the CCC plane ( $\mathrm{C}_{s}$ point group of symmetry). The two SPs named $\mathrm{SP}_{a a}$ are either the TSs of the rotation of the methylene related to C 2 , or of the rotation of the methylene of C3. The pathways to the two SPs are not totally symmetric, so that symmetry is broken at the VRI point named $\mathrm{VRI}_{a}$. A curious fact is that the $\mathrm{VRI}_{a}$ is reached with a further ring opening of the allyl radical (under a conrotatory motion of the terminal methylenes). Of course, the detection of the $\mathrm{VRI}_{a}$ without a guess in this direction is not possible. It could be caught, nevertheless, by a temporary use of gradient extremals [15]. They are implemented in the GAMESS-US [19, 20].

A full picture of the description by NTs of the allyl radical bowl is shown in Scheme 2. Note: the scheme fully plays on the lower surface of the minimum bowl. Near $\mathrm{SP}_{i 2}$ one could think a transition to $\mathrm{SP}_{d 1}$ on the upper surface: it is a small step in the coordinate space, as well as by a small


Scheme 2 Connection scheme of the allyl radical bowl. $\mathrm{M}_{a}$ : allyl radical minimum; the numbersindicate the order of H atoms, see Fig. 3 which here is the left, lower corner. Methylene 1 is the right group, methylene 2 the left group. $\mathrm{SP}_{a a}$ : TS of the rotation of one methylene group, thus allyl to allyl TS. $\mathrm{m}_{i} \perp$ : the corresponding methylene $i$ is quasi-orthogonal to the CCC plane. $\mathrm{SP}_{i 2}$ : SP of index 2 where both methylenes are quasi-orthogonal to the CCC plane. $\mathrm{VRI}_{a}$ : branching point of the singular NT between $\mathrm{M}_{a}$ and $\mathrm{SP}_{i 2}$, and the two different $\mathrm{SP}_{a a}$ 's (see Fig. 16 for the NT). Of course, there is also one $\mathrm{VRI}_{a}$ in the three other quadrants, correspondingly
energy jump. However, the tool of NTs is not well adapted to the task. Only pieces of NTs on the different surfaces are possible, compare Fig. 21 below. The schematic energy profile over all interesting points is given in Fig. 17.

## 9 The examination of steepest descent from $\mathbf{S P}_{c a}$

Last but not least, we report an SD calculation (without massweighting) through the interesting region after the $\mathrm{SP}_{c a}$ of the ring opening of cyclopropyl radical, see Figs. 18, 19 and 20. The SD curve runs after $\mathrm{SP}_{c a}$ the valley floor down and touches somehow the region of the $\mathrm{VRI}_{c a}$ point, compare Table A7 in the appendix where the SD node nearest to $\mathrm{VRI}_{c a}$ is reported. However, the dihedrals dih1 and dih3, as well as the bond length r 2 , do somehow deviate from the direct way to $\mathrm{VRI}_{c a}$. Though the singular NT after the $\mathrm{VRI}_{c a}$ is the crest line of the ridge, in contrast, the SD runs below the ridge region after $\mathrm{VRI}_{c a}$ downhill, still along a valley region, to the intersection where the transition between pi-radical to sigma-radical electronic state occurs. The aim is the allyl radical surface, and the transition is symbolized by stars *. The SD jumps on a ridge region on the lower surface, at $106^{\circ}$ of bond angle C 2 C 1 C 3 , and turns down along a flank, and later leaves the ridge region at the symbol ) (. There it arrives at


Fig. 17 Energies of the diverse points for the cyclopropyl radical to allyl radical ring opening
the convex minimum region. The finish is that it ends at the minimum $\mathrm{M}_{a}$ with all dihedrals and bond lengths at the corresponding values.

The $\mathrm{SP}_{a a}$ direction is nearly met by the dih2/-dih5 couple after the $\mathrm{VRI}_{c a}$. The single dihedral dih5 nearly touches the singular ridge NT belonging to $\mathrm{VRI}_{r l}$, in this projection. But the near consistence of dihedrals dih2 and -dih4 with the central branch of the singular NT of the lower $\mathrm{VRI}_{r l}$ is only an accident of this projection. However, dihedral dih3 again goes the other way, as well as dihedral dih1 and bond length r2.

The judgment to go down to the minimum $\mathrm{M}_{a}$, and not nearer to any version of the $\mathrm{SP}_{a a}$, is rendered in the region around $106^{\circ}$ of the bond angle C 2 C 1 C 3 , exactly at the electronic intersection. It becomes clear with the bond length coordinaterl in Fig. 19. An indicator for an early turn down of the SD bypassing the region of $\mathrm{SP}_{a a}$ is also the profile of the energy along the SD, see Fig. 20. It seems that around $106^{\circ}$ of the bond angle C 2 C 1 C 3 are two jumps in the energy, one shortly after the other, see again the star symbols. It could mean that there are two intersections, at least on the UHF PES.

The SD curve corresponds after the $\mathrm{SP}_{c a}$ for dihedrals dih2, dih3 and dih4 to a quasi-disrotatory character. However, the dihedral dih5 alone behaves conrotatory. After the intersection around $106^{\circ}$ of the bond angle C 2 C 1 C 3 , dih5 strongly changes the character and rotates exaggeratedly disrotatory. This is in accordance with the twofold intersection event [43]. But, the SD pathway shows no unique behavior, compare [44]. At all, the SD seems to be "navigated" by the singular NTs of the different VRI points. At the beginning, as one expects, the SD follows the valley into the region of $\mathrm{VRI}_{c a}$. And after the $\mathrm{VRI}_{c a}$, it goes down bordered by the different singular ridges. However, the SD does not feel the bumps of the PES like the corresponding branch of the singular NT through $\mathrm{VRI}_{c a}$. At


Fig. 18 Fat nodes indicate the steepest descent from $\mathrm{SP}_{c a}$ to allyl radical minimum. The asterisk symbol is the intersection to the PES of the allyl radical minimum, and the symbol) (means the border from ridge region to valley region near $120^{\circ}$ of bond angle C 2 C 1 C 3 . NTs

of thin points are the former systems of $\mathrm{VRI}_{c a}$ and $\mathrm{VRI}_{r l}$, for comparison, see Figs. 8 and 13. Note that the used projection sees the allyl and one version of $\mathrm{SP}_{a a}$ at the same place. Left coordinates dih2 and -dih5, right coordinates dih3 and -dih4


Fig. 19 Steepest descent from $\mathrm{SP}_{c a}$ to allyl radical minimum. See caption of Fig. 18. Left coordinates r 1 and r2, right coordinate d1


Fig. 20 left energy profile along the SD from $\mathrm{SP}_{c a}$. The two asterisks symbols two times indicate an electronic intersection of the UHF PES. The grid line is at the energy of the $\mathrm{SP}_{a a}$. Right example of an
the end, the final allyl radical minimum, $\mathrm{M}_{a}$, is connected to $\mathrm{M}_{c}$ and $\mathrm{SP}_{c a}$ by a disrotatory change [26]. Because of the intersection seam that crosses the descent region after the $\mathrm{VRI}_{c a}$, and which is represented by a steep precipice, the flow of SD lines from different beginnings (slightly "left" or "right" from VRI $_{c a}$, see also the next section) finds always the disrotatory allyl radical minimum. So to say, the $\mathrm{VRI}_{c a}$ point is relegated to uselessness by the lower electronic surface of the allyl radical side.

energy profile along two branches of the singular NT from $\mathrm{SP}_{c a}$ over $\mathrm{VRI}_{r}$ to $\mathrm{SP}_{d 1}$, see Figs. 12 or 21 . The peak emerges by the jump from one (not perfectly singular) branch to the next one

Technically, the SD calculation is done by simple Euler steps along the normalized contravariant gradient. The steplength used is 0.05 units (in the Bohr/radiant system). Because no further numerical effort is taken into account (compare, for example, [45]), there is some mild zigzagging, at the beginning in the SP region, as well as in the region of the minimum. However, the overall picture is not perturbed by this effect, especially because of the very small steplength.

## 10 The steepest descent from $\mathrm{VRI}_{r}$ and $\mathrm{SP}_{d 1}$

The corresponding SD curves are depicted in Fig. 21. The $\mathrm{VRI}_{r}$ is at the slope of the PES, and the SD from there goes down into the bowl of the minimum of the disrotatory allyl radical. At the other side, the two SD curves (and one for every dihedral in our double representation) run from $\mathrm{SP}_{d 1}$ in two opponent directions along the 1-dimensional valley floor down, and they find two different allyl radical minimums, either one of a disrotatory or one of a conrotatory structure.

The pathway of the SD in Fig. 21 is near the $\mathrm{SP}_{d 1}$ a rotation of one single methylene. The other methylene stays at the point and does not move. Thus, the SD is a pathway in direction to an $\mathrm{SP}_{a a}$. However, after $\approx 35^{\circ}$ of rotation, the intersection seam to the allyl radical surface is crossed. There the second methylene starts also a rotation and the radical goes down to the allyl minimum. In Fig. 21, left panel, the "upper" $\mathrm{SP}_{d 1}$ belongs to coordinate dih2, but the "lower" $\mathrm{SP}_{d 1}$ belongs to -dih5. The both col directions of the SP go along the pure dih2 (and not shown, a part of dih1 coordinate). A similar picture is the right panel, where dih3 is the active coordinate, but -dih4 stays. From point of view of the theory of NTs, there has to be a VRI point between $\mathrm{SP}_{d 1}$ and $\mathrm{SP}_{a a}$. We have done some search, but did not find it. The intersection seam again disturbs the pure theory of an NT-connection of the both SPs.

The $\mathrm{SP}_{d 1}$ is the outpost of the cyclopropyl radical side, of the upper surface. It is the "upper complement" to the $\mathrm{SP}_{i 2 a}$ on the allyl radical side, the $\sigma$-type radical electronic surface. One may speculate that there is also an $\mathrm{SP}_{i 2 u}$ on the upper surface near the two SPs of index one, $\mathrm{SP}_{d 1}$. Because, the relations of the two special points are different: For the methine bond, C 1 H 4 , the $\mathrm{SP}_{d 1}$ is a maximum at $180^{\circ}$, but the $\mathrm{SP}_{i 2 a}$ is a minimum. For the (quasi orthogonal to the C 2 C 1 C 3 plane) dihedrals of the two methylenes, the $\mathrm{SP}_{d 1}$ is a maximum for the rotation of one methylene, but a (very flat) minimum for a conrotatory
rotation of both methylenes, but the $\mathrm{SP}_{i 2 a}$ is the twofold maximum for the methylenes rotation. One eigenvector of the Hessian matrix with negative eigenvalue corresponds to the conrotatory motion while the second one is disrotatory. The direction of the SP valley of the $\mathrm{SP}_{d 1}$ is a linear combination of the H 4 C 1 bond vibration, and the rotation of only one methylene. So, the different index is explained. The points $\mathrm{SP}_{d 1}$ and $\mathrm{SP}_{i 2 a}$ are adjacent points, in coordinate space, as well as in their energy. But of course, between the two special points is the intersection seam. So, we can draw NTs between them, but they would cross the precipice. In Fig. 21, at the right-hand side in every panel, the NT is continued over the $\mathrm{SP}_{d 1}$ uphill, and then down the precipice to $\mathrm{SP}_{i 2}$. In contrast to this direction is the SD from $\mathrm{SP}_{d 1}$. It starts orthogonally in direction to an $\mathrm{SP}_{a a}$. The two symmetric forms of $\mathrm{SP}_{d 1}$ are characterized by a difference of only $3.3^{\circ}$ of the dihedrals of the two methylenes. However, this small difference is crucial for two very different directions of the corresponding SP cols in Fig. 21. From point of view of the theory, there has to be a further $\mathrm{SP}_{i 2}$ between the two $\mathrm{SPs}_{d 1}$ also on the upper surface. We have done some search, but we did not find it. The intersection seam disturbs this search.

The results are collected in Scheme 3. May be Scheme 3 and Fig. 21 can explain the trajectory results of Mann and Hase [27, 28].

In Fig. 22 we show the first two eigenvectors of the Hessian at $\mathrm{VRI}_{r}$. The direction of the eigenvector with zero eigenvalue points to the two SPs. The gradient leads to the structure where the unpaired electron is located in the plane C 2 C 1 C 3 corresponding to the electronic state ${ }^{2} \mathrm{~A}$ '. The $\mathrm{VRI}_{r}$ is at the slope of the PES, but the slope is small. Note that the singular branch to $\mathrm{SP}_{d 1}$ is a ridge line, but it does not need to be the crest line. We may assume that a dynamical trajectory can overfly the $\mathrm{VRI}_{r}$ region. When such a dynamical trajectory glides with high velocity over the $\mathrm{VRI}_{r}$ and runs further along, or near, the branch of the singular NT to $\mathrm{SP}_{d 1}$, then it has the possibility to slide



Fig. 21 Steepest descent (bullets) from $\mathrm{VRI}_{r}$ and $\mathrm{SP}_{d 1}$ to two versions of the allyl radical minimum. Thin points are the singular $\mathrm{NT}^{\text {of }} \mathrm{VRI}_{r}$. Left coordinates dih2 and -dih5, right coordinates dih3 and - dih4


Scheme 3 Possible flow of singular NTs (thick arrows) and some SD lines (thin arrows) for the ring opening after $\mathrm{SP}_{c a}$. (For symbols see the Captions of former schemes.) Asterisk depicts an electronic intersection. Corresponding arrows for NTs are drawn only symbolically. In contrast, SD lines can jump down and continue after the intersection seam
down at the "conrotatory" side of the flat ridge which itself leads to $\mathrm{SP}_{d 1}$. So to say, the branch of the singular NT from $\mathrm{VRI}_{r}$ to $\mathrm{SP}_{d 1}$ is a kind of a TS demarcation line [46]. If it is overcome then the conrotatory side of the PES is possible. The $\mathrm{VRI}_{r}$ does not be an obstacle. It is only a little higher in energy than the $\mathrm{SP}_{c a}$, and the ridge from $\mathrm{VRI}_{r}$ goes very slowly down to $\mathrm{SP}_{d 1}$, see the energy profile in the right panel of Fig. 20. The ridge to $\mathrm{SP}_{d 1}$ is throughout on the upper electronic PES. A hypothetical, dynamical trajectory could connect the $\mathrm{SP}_{c a}$ with the minimum of the conrotatory structure of the allyl radical, being a result of Mann and Hase [27, 28].

A qualitative similar PES situation was found for another reaction system
$\mathrm{CH}_{3} \mathrm{O} \leftrightarrow \mathrm{H}_{2}+\mathrm{HCO} \leftrightarrow \mathrm{H}+\mathrm{H}_{2} \mathrm{CO}$.
A two-dimensional section of the corresponding PES is depicted in Fig. 4 of [47]. The reaction was named a reaction over a barrier but not through an SP. Further examples of reactions, which do not follow the post-TS IRC, are reported [48-52]. Of course, a reaction which
does not follow the post-TS IRC is not so easy to explore with the tool of SD curves [47]. It could be explored better if one uses the tool of NTs.

## 11 Discussion

The PES of the ring opening of the cyclopropyl radical is notoriously curvilinear. To see anything at all, we concentrate our interest to projections to the 5 dihedrals, and the two distances of the C -atoms of the 8 -atomic molecule. It may be problematic to treat only these projections: there are still 5 distances and 5 angles of the H atoms not treated here (except in the reported tables in the appendix).

We have tested the reliability of the use of the UHF PES by calculation of the corresponding CASSCF energies of representative points of the UHF PES. The results in Table 2 show an energy difference that is constant for nearly all the points, meaning that UHF underestimates the energy about $8-9 \mathrm{kcal} / \mathrm{mol}$. Only for some points being near to the electronic intersection this energy difference raises above $10 \mathrm{kcal} / \mathrm{mol}$. This demonstrates, that UHF is a good approximation, as there is a constant error, and not an error depending on the current point. Curiously, the UHF PES around the symmetric $\mathrm{SP}_{i 2}$ of the allyl bowl has a strange energy difference, as well as the quasi-symmetric $\mathrm{SP}_{d 1}$. The energy calculated at the UHF level is lower than that of the CAS calculation, but this, although unusual, is possible mathematically. The reason is that here the quadruplet electronic state (occupation $1.0,1.0,1.0$ ) using UHF wave function for the energy is lower with respect to the doublet state and CASSCF. It means that in this region the quadruplet electronic state is the lowest electronic state rather than the doublet state. In others words, in this region the calculation of CASSCF corresponds to an excited state.

## 12 Conclusion

With this work, we open the possibility to channelize the ring opening of the cyclopropyl radical between some

Fig. 22 left zero eigenvector at $\mathrm{VRI}_{r}$. It points along the two branches from $\mathrm{SP}_{c a}$ to $\mathrm{SP}_{d 1}$ and marks the level line of curvature zero. Right second eigenvector in direction of the gradient at $\mathrm{VRI}_{r}$. It points along the other two branches


Table 2 Comparison of UHF and CAS energy values of representative UHF points with basis set 6-31 G(d)

| Point | Energy <br> UHF au | Energy <br> CAS au | Difference <br> acal/mol |
| :--- | :--- | :--- | :---: |
| $\mathbf{M}_{c}^{\mathrm{b}}$ | -116.41492 | -116.42759 | 7.95 |
| $\mathrm{VRI}_{c}$ | -116.37773 | -116.39038 | 7.94 |
| $\mathrm{SP}_{H 4}$ | -116.40935 | -116.42106 | 7.35 |
| $\mathrm{SP}_{i 2 c}$ | -116.34794 | -116.36014 | 7.66 |
| $\mathrm{SP}_{c a}$ | -116.37268 | -116.39261 | 12.51 |
| $\mathrm{VRI}_{c a}$ | -116.37565 | -116.39027 | 9.17 |
| $\mathrm{VRI}_{r}$ | -116.37204 | -116.38154 | 5.96 |
| $\mathrm{VRI}_{s}$ | -116.30119 | -116.31365 | 7.82 |
| $\mathrm{IRC}_{1}$ | -116.37819 | -116.39204 | 8.70 |
| $\mathrm{IRC}_{2}$ | -116.39950 | -116.42460 | 15.75 |
| $\mathrm{SP}_{\mathrm{d} 1}^{c}$ | -116.37394 | -116.36720 | -4.23 |
| $\mathrm{SP}_{i 2 a}$ | -116.37782 | -116.35932 | -11.61 |
| $\mathrm{IRC}_{3}^{d}$ | -116.41226 | -116.43681 | 15.41 |
| $\mathrm{VRI}_{r 11}$ | -116.42231 | -116.44542 | 14.50 |
| $\mathrm{SP}_{a a}$ | -116.43851 | -116.46178 | 14.60 |
| $\mathrm{VRI}_{a}$ | -116.44217 | -116.45563 | 8.45 |
| $\mathrm{M}_{a}$ | -116.46810 | -116.48358 | 9.71 |

${ }^{\text {a }}$ Factor: $1 \mathrm{au}=627.509 \mathrm{kcal} \mathrm{mol}^{-1}$
b The upper table concerns the PES of the cyclopropyl radical side
${ }^{\text {c }}$ See text for the negative differences
${ }^{\text {d }}$ The lower table concerns the PES of the allyl radical side
singular NTs. The known SP of the ring opening, $\mathrm{SP}_{c a}$ [25], is now included in a row of VRI points, as well as SPs of index 2, on the UHF PES. The tool to find these special points is a variational application of Newton trajectories (NT) proposed in [18].

The breakdown of the symmetry of the cyclopropyl radical, or the allyl radical takes place at the pre-transition state VRI point, $\mathrm{VRI}_{c}$, and at point $\mathrm{VRI}_{a}$. Up to these two points, a possible reaction path coming from the minimums could be thought to be fully symmetric. But the transition states $\mathrm{SP}_{c a}$ and $\mathrm{SP}_{a a}$ are not symmetric, and so the branches leading to these SPs have to be not symmetric as well. The corresponding archetypes are Fig. 4 left panel, and Fig. 16 right panel. Of course, any "direct" pathway between minimum $\mathrm{M}_{c}$ and $\mathrm{SP}_{c a}$, for example, can be a model of the reaction path including the SD from $\mathrm{SP}_{c a}$. The singular NT through the $\mathrm{VRI}_{c}$ only marks a border for such RP models marking a so-called reaction channel [53-55].

After the $\mathrm{SP}_{c a}$ of the ring opening, and before the adjacent $\mathrm{SP}_{a a}$ in the allyl bowl, has to be at least one further VRI point. We propose to assume the $\mathrm{VRI}_{c a}$ for that event. However, the case of electronic intersections of different PESs disturbs a usual application for this molecule of the NT theory concerning the existence of VRI points in certain relations between the stationary points
[8, 9]. Rising uphill from $\mathrm{SP}_{a a}$ or $\mathrm{SP}_{i 2 a}$, correspondingly, we found further "systems" of VRI points of rpVRI character of higher index. Their singular NTs may not directly lead to $\mathrm{SP}_{c a}$, and they are lost somewhere in the PES mountains of non-chemical high energies. However, they narrow the region where one can guess the pathway of the ring opening. This is demonstrated by a SD calculation, at least, which finds its way through the different singular NTs. These set the limits for the SD form $\mathrm{SP}_{c a}$ to $\mathrm{M}_{a}$.

On the other hand, the use of SD separates the different interesting points of the PES (SPs, VRI points) into different regions of attraction of the dis or conrotatory minimums. The ring opening shows a very mixed character of con and disrotatory motion of the methylenes if one traces the different singular NTs, as well as the SD from $\mathrm{SP}_{c a}$. Thus, the question of the character of the ring opening is that neither a clear conrotatory nor a clear disrotatory motion takes place [44]. The summary result of NTs or SD from $\mathrm{SP}_{c a}$ down the SP valley is the disrotatory minimum. However, the singular NT along a ridge region from $\mathrm{SP}_{c a}$ to $\mathrm{SP}_{d 1}$ could explain a possibility for a conrotatory ring opening. Thus, the diverse singular NTs set the limits for usual static reaction pathways. But they may open the insight for results of dynamical trajectories, compare [27,28], where some trajectories go to a conrotatory ring opening.

A side result in this work is the assignment of a so-called mixed character VRI point on the PES of a real molecule, the $\mathrm{VRI}_{H 4}$. The kind of VRI point was recently described [15, 17]. A further side result in this work is the detection, and possible meaning of so-called bumps on the PES. They are indicated by double bends of the NTs, see Fig. 10.

With the reaction of the cyclopropyl radical ring opening to allyl radical, we have shown that the NT methodology provides a tool to explore the topology of the potential energy surface. The so-called singular Newton trajectories acquire special importance. These curves pass through the VRI points and the corresponding analysis of each branch that crosses the VRI provides information on the structure and form of the surface around this point. They are related with the bifurcation of valleys. This work also seeks to pave the way to further fruitful applications of the mathematics of VRI points and the bifurcation of NTs in other molecules.

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